

Handout of more complete proof of
 $\mu_i = \mu_i^o + RT \ln P_i$ for mixture of ideal gases

The object is a complete derivation of the relationship between chemical potential μ_i and partial pressure by $\mu_i = \mu_i^o + RT \ln P_i$. This relation and its counterparts for concentration of solutes is so important that it deserves a **HANDOUT**:

We calculate ΔG_{mixing} for the same arrangement used previously to calculate ΔS_{mixing} :



For ideal gas

$$\Delta H_{\text{mixing}} = 0; \quad \Delta S_{\text{mixing}} = -(n_A + n_B)R[X_A \ln X_A + X_B \ln X_B] = -R[n_A \ln X_A + n_B \ln X_B]$$

(1.1)

so

$$(\Delta G_{\text{mixing}})_{T,P} = -T\Delta S_{\text{mixing}} = RT[n_A \ln X_A + n_B \ln X_B]$$

also

$$\Delta G_{\text{mix}} = G_{A+B} - (G_A + G_B)$$

(1.2)

with

$$G_{A+B} = n_A \mu_A + n_B \mu_B$$

$$G_A = G_A^o + n_A RT \ln \left(\frac{P}{1 \text{ atm}} \right)$$

(1.3)

$$G_B = G_B^o + n_B RT \ln \left(\frac{P}{1 \text{ atm}} \right)$$

with $G_A^o = n_A \bar{G}_A^o$ $G_B^o = n_B \bar{G}_B^o$

thus

$$\Delta G_{\text{mix}} = n_A \mu_A + n_B \mu_B - n_A \bar{G}_A^o - n_A RT \ln \left(\frac{P}{1 \text{ atm}} \right) - n_B \bar{G}_B^o - n_B RT \ln \left(\frac{P}{1 \text{ atm}} \right)$$

(1.4)

and

$$\Delta G_{\text{mix}} = RT[n_A \ln X_A + n_B \ln X_B]$$

equating and rearranging

$$n_A \left[\mu_A - \bar{G}_A^o - RT \ln \left(\frac{P}{1 \text{ atm}} \right) - RT \ln X_A \right] + n_B \left[\mu_B - \bar{G}_B^o - RT \ln \left(\frac{P}{1 \text{ atm}} \right) - RT \ln X_B \right] = 0$$

(1.5)

next we differentiate both side by $\left(\frac{\partial}{\partial n_A} \right)_{n_B, T, P}$

$$\left[\mu_A - \bar{G}_A^o - RT \ln \left(\frac{P}{1 \text{ atm}} \right) - RT \ln X_A \right] + n_A \left(\frac{\partial \mu_A}{\partial n_A} \right)_{n_B, T, P} - n_A RT \frac{1}{X_A} \left(\frac{\partial X_A}{\partial n_A} \right)_{n_B, T, P}$$

(1.6)

$$+ n_B \left(\frac{\partial \mu_B}{\partial n_A} \right)_{n_B, T, P} - n_B RT \frac{1}{X_B} \left(\frac{\partial X_B}{\partial n_A} \right)_{n_B, T, P} = 0$$

to get the derivatives of X_A and X_B with respect to n_A :

$$X_A = \frac{n_A}{n_A + n_B} \quad X_B = \frac{n_B}{n_A + n_B}$$

so

$$\left(\frac{\partial X_A}{\partial n_A} \right)_{n_B, T, P} = \frac{1}{n_A + n_B} - \frac{n_A}{(n_A + n_B)^2} = \frac{n_B}{(n_A + n_B)^2} = \frac{X_A X_B}{n_A} \quad (1.7)$$

and

$$\left(\frac{\partial X_B}{\partial n_A} \right)_{n_B, T, P} = -\frac{n_B}{(n_A + n_B)^2} = -\frac{X_B X_B}{n_B} = -\frac{X_B^2}{n_B}$$

substituting (1.7) into (1.6)

$$\left[\mu_A - \bar{G}_A^\circ - RT \ln \left(\frac{P}{1 \text{ atm}} \right) - RT \ln X_A \right] + n_A \left(\frac{\partial \mu_A}{\partial n_A} \right)_{n_B, T, P} - n_A RT \frac{1}{X_A} \frac{X_A X_B}{n_A} + n_B \left(\frac{\partial \mu_B}{\partial n_A} \right)_{n_B, T, P} - n_B RT \frac{1}{X_B} \left(\frac{-X_B^2}{n_B} \right) = 0 \quad (1.8)$$

note that the last four terms involve two cancellations:

$$-n_A RT \frac{1}{X_A} \frac{X_A X_B}{n_A} - n_B RT \frac{1}{X_B} \left(\frac{-X_B^2}{n_B} \right) = 0 \quad \text{by algebra} \quad (1.9)$$

and

$$n_A \left(\frac{\partial \mu_A}{\partial n_A} \right)_{n_B, T, P} + n_B \left(\frac{\partial \mu_B}{\partial n_A} \right)_{n_B, T, P} = 0 \quad \text{by the Gibbs - Duhem relationship proved elsewhere}$$

FINALLY we are left with:

$$\left[\mu_A - \bar{G}_A^\circ - RT \ln \left(\frac{P}{1 \text{ atm}} \right) - RT \ln X_A \right] = 0 \quad (1.10)$$

$$\left[\mu_A = \bar{G}_A^\circ + RT \ln \left(\frac{P}{1 \text{ atm}} \right) + RT \ln X_A \right]$$

ALMOST THERE

$$\mu_A = \bar{G}_A^\circ + RT \ln \left(\frac{X_A P}{1 \text{ atm}} \right) \quad (1.11)$$

with $P X_A = \frac{n_A}{n_A + n_B} P = P_A$ (partial pressure of gas A)

and $\mu_A^\circ(T) \equiv \mu_A(T, 1 \text{ atm, pure A}) = \bar{G}_A^\circ(T)$

WE GET from (1.10)

$$\boxed{\mu_A(T) = \mu_A^\circ(T) + RT \ln \left(\frac{P_A}{1 \text{ atm}} \right)} \quad (1.12)$$

HOORAY !!