

Chemistry 163B Winter 2012
Handouts for Third Law and Entropy of Mixing
(ideal gas, distinguishable molecules)

Chemistry 163B
Absolute Entropies
and
Entropy of Mixing

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APPENDIX A: ΔH_f° , ΔG_f° , BUT S° (no Δ , no "sub f")


Substance	ΔH_f° ΔH_f° (kJ mol ⁻¹)	ΔG_f° ΔG_f° (kJ mol ⁻¹)	S° S° (J mol ⁻¹ K ⁻¹)	$C_{p,m}$ $C_{p,m}$ (J mol ⁻¹ K ⁻¹)	Atomic or Molecular Weight (amu)
Carbon					
Graphite(s)	0	0	5.74	8.52	12.011
Diamond(s)	1.89	2.90	2.38	6.12	12.011
C(s)	-716.7	-671.2	158.1	20.8	12.011
CO(g)	-110.5	-137.2	197.7	29.1	28.011
Hydrogen					
H ₂ (g)	0	0	130.7	28.8	2.016
H ₂ O(g)	-241.8	-228.6	188.8	33.6	18.015
H ₂ O(l)	-285.8	-237.1	70.0	75.3	18.015
H ₂ O(s)	-286.3	-236.6	48.0	36.2 (273 K)	18.015
H ₂ O(l)	-136.3	-105.6	232.7	43.1	34.015
H ⁺ (aq)	0	0	0	1.008	1.008
OH ⁻ (aq)	-230.0	-157.24	-33.9	17.01	17.01
Oxygen					
O ₂ (g)	0	0	205.2	29.4	31.999
O ₃ (g)	247.2	231.7	161.1	21.9	15.999
O ₂ (g)	142.7	163.2	238.9	39.2	47.998
OH ₂ (g)	39.0	34.22	183.7	29.9	17.01
OH ⁻ (aq)	-230.0	-157.2	-10.9	17.01	17.01

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Third Law of Thermodynamics

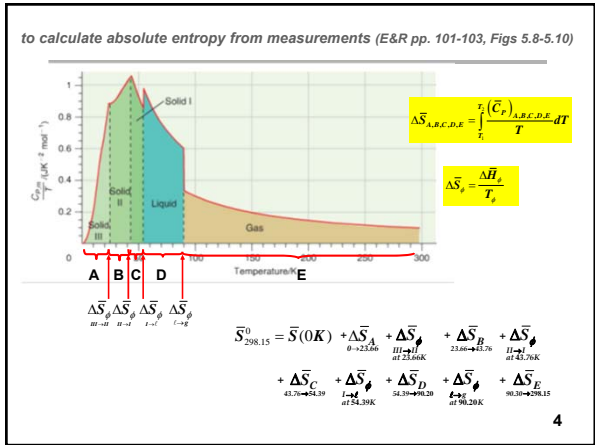
The entropy of any perfect crystalline substance approaches 0 as T → 0K

S = k ln W



for perfectly ordered crystalline substance
W → 1 as T → 0K ⇒ S → 0

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full calculation of S_{298}° for O₂(g) (Example Problem 5.9, E&R pp103-104 [96-97]...)

	$\Delta\bar{S}^\circ$ J K ⁻¹ mol ⁻¹
$\bar{S}^\circ(0K)$	0
$\Delta\bar{S}_A^\circ(0 \rightarrow 23.66)$	8.182
$\Delta\bar{S}_B^\circ(III \rightarrow II \text{ at } 23.66K)$	3.964
$\Delta\bar{S}_C^\circ(23.66 \rightarrow 43.76)$	19.61
$\Delta\bar{S}_D^\circ(II \rightarrow I \text{ at } 43.76K)$	16.98
$\Delta\bar{S}_E^\circ(43.76 \rightarrow 54.39)$	10.13
$\Delta\bar{S}_F^\circ(I \rightarrow \ell \text{ at } 54.39K)$	8.181
$\Delta\bar{S}_G^\circ(54.39 \rightarrow 90.20)$	27.06
$\Delta\bar{S}_H^\circ(\ell \rightarrow g \text{ at } 90.20K)$	75.59
$\Delta\bar{S}_I^\circ(90.20 \rightarrow 298.15)$	35.27
total	204.9 J K⁻¹ mol⁻¹

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$\Delta S_{\text{reaction}}$ from absolute entropies

$n_A A + n_B B \rightarrow n_C C + n_D D$ at 298K

$\Delta S_{\text{reaction}} = n_C (\bar{S}_{298}^\circ)_C + n_D (\bar{S}_{298}^\circ)_D - n_A (\bar{S}_{298}^\circ)_A - n_B (\bar{S}_{298}^\circ)_B$

$\Delta S_{\text{reaction}}^\circ(298K) = \sum_i \nu_i (\bar{S}_{298}^\circ)_i$

$(\bar{S}_{298}^\circ)_i$ are 3rd Law entropies (e.g. Appendix A)

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qualitative factors affecting molecular entropy

- Higher T $\Rightarrow \left(\frac{\partial S}{\partial T}\right)_P = \frac{C_p}{T} > 0$
- Higher P $\Rightarrow \left(\frac{\partial S}{\partial P}\right)_T = -\left(\frac{\partial V}{\partial T}\right)_P < 0$
usually
- Phase S(g) vs S(l) vs S(s)
(in a reaction the side with the greater number of moles of gas generally has higher S)
- Mixing or dissolving of components
(l+l), (s+s), (l+s), (g+g) solutions
 \Rightarrow
- (g + l) or (g + s) solution \Rightarrow

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more qualitative factors affecting molecular entropy

- substances with higher mass have

S_{298}^0	$F_2(g)$	$Cl_2(g)$	$Br_2(g)$	$I_2(g)$
	202.78	223.07	245.46	260.69

(more closely spaced rotational and vibrational levels)
- more rigid substances have

S_{298}^0	C(gr)	C(dia)
	5.74	2.377

J K⁻¹ mol⁻¹
- more complex substances have

S_{298}^0	HF (g)	H ₂ O (g)	D ₂ O (g)	
	20	18	20	amu
	173.78	188.83	198.34	J K ⁻¹ mol ⁻¹

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calculating entropy (see summary on review handout)

- Thermal properties of entropy and entropy calculations
- ✓ $dS = \frac{dq_{rev}}{T}$; $\Delta S = \int \frac{dq_{rev}}{T}$; $\int \frac{dq_{rev}}{T} = 0$
- ✓ $\Delta S \geq \int \frac{dq}{T}$; $0 \geq \int \frac{dq}{T}$; (*= for reversible process; > for spontaneous [“real”] process*)
- ✓ $\Delta S_{universe} = \Delta S_{system} + \Delta S_{surroundings} \geq 0$
- ✓ S is a state function; dS is an exact differential
- Dependence of S on
 - T: $\left(\frac{\partial S}{\partial T}\right)_P = \frac{C_p}{T}$; $\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_v}{T}$
 - P: $\left(\frac{\partial S}{\partial P}\right)_T = -\left(\frac{\partial V}{\partial T}\right)_P$
 - V: $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_P$
- Phase: $\Delta S = \frac{\Delta H_{transition}}{T_{transition}}$
- Calculation of entropy changes for changes in P, V, T, phase
- Third Law and calculations using Third Law Entropies: S^o(T)
- ✓ $\Delta S_{mix}^o(T) = -R \sum_i \nu_i \ln \nu_i$
- Entropy of mixing: $\Delta S = -n_{total} R \sum_i X_i \ln X_i$, where $X_i = \frac{n_i}{n_{total}}$

Lecture 9, slide 3

from lecture on 2nd Law and probability (abstract)

- Discrete: W_{tot} did not change during an arbitrary reversible expansion (W_{tot} = 2^N)
- Discrete: W_{tot} increased in isothermal reversible expansion (W_{tot} = 2^N)
- Discrete: W_{tot} decreased with T increase (2^N)
- Discrete: W_{tot} decreased with T decrease (2^N)
- All T > 0, W > 1

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HW#5 Prob 31

31.

- a. Prove, in general (but $g_{surroundings} = 0$, $ds = 0$), the relationship in equation 3.19 E&R:

$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial P}{\partial T}\right)_V - P$
- b. Evaluate $\left(\frac{\partial U}{\partial V}\right)_T$ for an ideal gas.

Sample midterm 3a
- c. Evaluate $\left(\frac{\partial U}{\partial T}\right)_V$ for a Van der Waals gas $\left(P + \frac{a}{V^2}\right)(V-b) = RT$

analogous to 31a
- d. Interpret your (correct) results for parts (b) and (c).
- e. Prove, in general (but $g_{surroundings} = 0$, $ds = 0$), the relationship in equation 3.44 E&R:

$\left(\frac{\partial U}{\partial P}\right)_T = T \left(\frac{\partial V}{\partial T}\right)_P - V$

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the relationships

<p>definitions:</p> <p>U \equiv internal energy</p> <p>H \equiv U + PV</p> <p>A \equiv U - TS</p> <p>G \equiv H - TS</p> <p style="margin-top: 10px;">heat and temperature:</p> <p>$dq_v = n \bar{C}_v dT$ $dq_p = n \bar{C}_p dT$</p> <p>$dU = dq + dw = dq - PdV$</p> <p>$dS = \frac{dq_{rev}}{T}$ $dq = Tds$</p> <p>$\left(\frac{\partial S}{\partial T}\right)_V = \frac{n \bar{C}_v}{T}$ $\left(\frac{\partial S}{\partial T}\right)_P = \frac{n \bar{C}_p}{T}$</p>	<p>differentials of state functions:</p> <p>$dU = TdS - PdV$</p> <p>$dH = TdS + VdP$</p> <p>$dA = -SdT - PdV$</p> <p>$dG = -SdT + VdP$</p>
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do some examples:

HW#5 Prob 31a:
derive E&R equation 3.19 'LATER is NOW'

$$\left(\frac{\partial U}{\partial V}\right)_T = ???$$

in terms of P, V, T and their derivatives

technique applies to HW#6 Prob: 31e

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do another example:

One mole of $\text{CO}_2(\text{g})$ is expanded isothermally and reversibly from V_1 to V_2 . Using the van der Waals equation of state

$$\left(P + \frac{a}{\bar{V}^2}\right)(\bar{V} - b) = RT$$

to describe $\text{CO}_2(\text{g})$ calculate w , ΔU , q , and ΔS in terms of V_1 and V_2 and the van der Waals constants a and b .

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Entropy of Mixing of Ideal Gasses

(EXTRA)

E&R \approx Sec 6.6

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Entropy of mixing for ideal gas (distinguishable particles)

- isolated from surroundings
 $q_{\text{sys}} = q_{\text{surr}} = 0$
 $w = 0$
- $\Delta S_{\text{surr}} = 0$
- $\Delta S_{\text{universe}} > 0$
- $\Delta S_{\text{sys}} > 0 = ???$

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reversible isothermal expansion to total volume $\Delta S = nR \ln \frac{V_{\text{final}}}{V_{\text{initial}}}$

$\Delta S = n_a R \ln \frac{V_a + V_b}{V_a}$ $\Delta S = n_b R \ln \frac{V_a + V_b}{V_b}$

n_a
 $V_a + V_b$
 T
 $P_a = n_a \frac{RT}{V_a + V_b}$

n_b
 $V_a + V_b$
 T
 $P_b = n_b \frac{RT}{V_a + V_b}$

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reversible recombining of volumes

permeable to b
not permeable to a

permeable to a
not permeable to b

- T const
- $\Delta U = 0$
- $P_{\text{ext}} = 0 \rightarrow w = 0$
- $q_{\text{rev}} = 0$
- $\Delta S > 0$

- T const
- $\Delta U = 0$
- $P_{\text{ext}} = 0 \rightarrow w = 0$
- $q_{\text{rev}} = 0$
- $\Delta S = 0$

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some arithmetic

note 'flip' leads to $-\ln$ note 'flip' leads to $-\ln$

$$\Delta S_{\text{total}} = n_a R \ln \frac{V_a + V_b}{V_a} + n_b R \ln \frac{V_a + V_b}{V_b}$$

(only isothermal expansion steps contribute)

partial pressure: $P_i = \frac{n_i}{n_{\text{total}}} P_{\text{total}}$ mole fraction: $X_i = \frac{n_i}{n_{\text{total}}}$

ideal gas: at start $P V_a = n_a R T$; $P V_b = n_b R T$ divide by

at end $P(V_a + V_b) = (n_a + n_b) R T$ ←

$$\frac{V_a}{V_a + V_b} \rightarrow \frac{n_a}{n_a + n_b} = X_a$$

$$\frac{V_b}{V_a + V_b} \rightarrow \frac{n_b}{n_a + n_b} = X_b$$

$\Delta S = -n_a R \ln X_a - n_b R \ln X_b$

and per mole

$$\frac{\Delta S}{n_{\text{total}}} = \Delta \bar{S} = -X_a R \ln X_a - X_b R \ln X_b$$

$$\Delta \bar{S}_{\text{mixing}} = -R \sum_i X_i \ln X_i$$

$$\Delta S_{\text{mixing}} = -n_{\text{total}} R \sum_i X_i \ln X_i$$

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calculating entropy (see summary on review handout) \Rightarrow

- Thermal properties of entropy and entropy calculations
- ✓ $ds = \frac{\delta q_{rev}}{T}$; $\Delta S = \int \frac{\delta q_{rev}}{T}$; $\oint \frac{\delta q_{rev}}{T} = 0$
- ✓ $\Delta S \geq \int \frac{\delta q}{T}$; $0 \geq \oint \frac{\delta q}{T}$; (= for reversible process; > for spontaneous [real] process)
- ✓ $\Delta S_{universe} = \Delta S_{system} + \Delta S_{surroundings} \geq 0$
- ✓ S is a state function; ds is an exact differential
- ✓ Dependence of S on
 - T: $\left(\frac{\partial S}{\partial T}\right)_P = \frac{C_p}{T}$; $\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_v}{T}$
 - P: $\left(\frac{\partial S}{\partial P}\right)_T = -\left(\frac{\partial V}{\partial T}\right)_P$
 - V: $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$
- ✓ Phase $\Delta S = \frac{\Delta H_{equilibrium phase change}}{T_{equilibrium phase change}}$
- ✓ Calculation of entropy changes for changes in P, V, T, phase
- ✓ Third Law and calculations using Third Law Entropies: $S^\circ(T)$
- ✓ $\Delta S^\circ_{mix}(T) = -R \sum_i x_i \ln x_i$
- ✓ Entropy of mixing: $\Delta S = -n_{total} R \sum_i x_i \ln x_i$, where $x_i = \frac{n_i}{n_{total}}$

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What formulas I have to memorize for midterm and final ?

Handout for CHEMISTRY 163B
ENTROPY for C, H, N, O and Cl based on a given mass

Use differentials for ΔS , ΔG , and ΔG based on the G as a function

For variable volume path, ideal gas

- $P_1 V_1 = P_2 V_2$ (ideal gas, variable path)
- $\int \frac{1}{V} dV = \ln V$ (ideal gas, variable path, $P \neq \text{constant}$)
- $\frac{d}{dx} \ln x = \frac{1}{x}$ (ideal gas, variable path)

Change in entropy

- $\Delta S = nR \ln \frac{V_2}{V_1} = nR \ln \frac{P_1}{P_2}$ (ideal gas, variable path)

For real gas

- $\Delta S = nR \ln \frac{V_2}{V_1} + \int_{P_1}^{P_2} \left(\frac{\partial S}{\partial P}\right)_T dP$
- $\left(\frac{\partial S}{\partial P}\right)_T = -\left(\frac{\partial V}{\partial T}\right)_P$ (Maxwell relation)
- $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$ (Maxwell relation)

For ideal gas, variable path

- $\Delta S = nR \ln \frac{V_2}{V_1} + nC_p \ln \frac{T_2}{T_1} - nR \ln \frac{P_2}{P_1}$

Use the following relations

- $\left(\frac{\partial S}{\partial T}\right)_P = \frac{C_p}{T}$
- $\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_v}{T}$
- $\left(\frac{\partial S}{\partial P}\right)_T = -\left(\frac{\partial V}{\partial T}\right)_P$
- $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$


(see handout for more derivations)

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Gibb's Paradox: distinguishable vs identical (don't fret)

n_a
 V_a
T, P

n_b
 V_b
T, P



n_a
 V_a
T, P

n_b
 V_b
T, P

$n_a = n_b$
 $V_a = V_b$
 $\Delta S_{total} = 2n_a R \ln 2$

$2n_a$
 $2V_a$
T, P

$\Delta S_{total} = 0$

$\Delta S_{total} = n_a R \ln \frac{V_a + V_b}{V_a} + n_b R \ln \frac{V_a + V_b}{V_b}$

CORRECT **CORRECT TOO**

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Gibb's Paradox: distinguishable vs identical (don't fret)

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Full Paper

Gibb's Paradox and the Definition of Entropy

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The Gibbs Paradox and the Distinguishability of Identical Particles

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statistics, Chemistry 163C

Some references:


www.mdpi.org/lin/entropy/Chua-Thermodynamics.ppt

http://baves.wustl.edu/eti/articles/gibbs_paradox.pdf

http://www.nyu.edu/classes/tuckerman/stat_mech/lectures/lecture_6/node5.html

<http://www.istor.org/stable/186086>


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where we have been

and

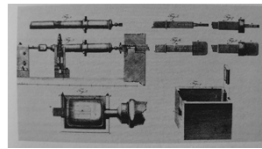
what's next



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observations: thermo = heat

- Count Rumford, 1799
- observed water turning into steam when canon barrel was bored
- work \leftrightarrow heat






FIGURE 11. An illustration from Rumford's paper, "An Inquiry Concerning the Source of the Heat Which is Evolved by Friction," showing the apparatus used by him in the cannon boring experiment. Figure 1, upper left, shows the cannon as received from the foundry, and Figure 2, below, shows it mounted in the machine used for boring. (Reproduced with the permission of Harvard University Press.)


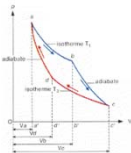

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observations: mechanical efficiency of steam engine

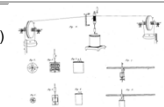


- Sadi Carnot, 1824
- efficiency of engines

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guiding principles

- Conservation of heat and work (Joule, 1845)
1st LAW OF THERMODYNAMICS
- Clausius, 1860 Entropy
2nd LAW OF THERMODYNAMICS
- Boltzmann, late 19th century,
molecular picture of entropy


**the thermodynamic functions U, H, and S
(1st and 2nd laws)**

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"Applications"

How does knowledge about efficiencies of steam engines, mechanical systems, etc. relate to processes in chemical, biological, and geological systems?

ANSWERED BY:



J. W. Gibbs- arguably the first great American scientist who combined the concepts of heat and entropy and proposed "[Gibbs] Free Energy", G, a thermodynamic state function that leads to a whole spectrum of applications

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End of Lecture

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In 1872, long before women were admitted to Austrian universities, he met Henriette von Aigentler, an aspiring teacher of mathematics and physics in Graz. She was refused permission to audit lectures, unofficially Boltzmann advised her to appeal, which she did successfully.

In Vienna, Boltzmann not only taught physics but also lectured on philosophy. Boltzmann's lectures on natural philosophy were very popular, and received a considerable attention at that time. His first lecture was an enormous success. Even though the largest lecture hall had been chosen for it, the people stood all the way down the staircase. Because of the great successes of Boltzmann's philosophical lectures, the Emperor invited him for a reception at the Palace.

On September 5, 1906, while on a summer vacation in Duino, near Trieste, Boltzmann hung himself during an attack of depression.^[19] He is buried in the Viennese Zentralfriedhof; his tombstone bears the inscription

$S = k \cdot \log W$

↑



Ludwig Boltzmann

To quote Planck, "The logarithmic connection between entropy and probability was first stated by L. Boltzmann in his kinetic theory of gases".^[10] This famous formula for entropy S is^{[11][12]}

$S = k \log_e W$

where $k = 1.3806505(24) \times 10^{-23} \text{ J K}^{-1}$ is Boltzmann's constant, and the logarithm is taken to the natural base e. W is the *Wahrscheinlichkeit*, the frequency of occurrence of a macrostate^[13] or, more precisely, the number of possible microstates corresponding to the macroscopic state of a system — number of