

Chemistry 163B Winter 2020

Lecture 9- Carnot Arithmetic and Efficiency

Chemistry 163B
Lecture 09


Carnot Arithmetic

Challenged Penmanship

Notes

[see handout: Carnot Arithmetic](#) →

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roadmap for second law 

1. Phenomenological statements (what is ALWAYS observed)
2. Ideal gas Carnot [reversible] cycle efficiency of heat → work (Carnot cycle transfers heat only at T_U and T_L)
3. Any cyclic engine operating between T_U and T_L must have an equal or lower efficiency than Carnot OR VIOLATE one of the phenomenological statements (observations)
4. Generalize Carnot to any reversible cycle (E&R fig 5.4)
5. Show that for this REVERSIBLE cycle
 $q_U + q_L \neq 0$ (dq inexact differential $\oint dq \neq 0$)
but
 $\frac{q_U}{T_U} + \frac{q_L}{T_L} = 0$ (something VERY VERY special about $\frac{dq_{rev}}{T}$; $\oint \frac{dq_{rev}}{T} = !!!$)
6. S, entropy and spontaneous changes

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from lecture on 2nd Law and probability (disorder)

- Disorder, **W**, did not change during an adiabatic reversible expansion ($q_{rev} = 0$)
- Disorder, **W**, increased in isothermal reversible expansion ($q_{rev} > 0$)
- Disorder, **W**, increased with T increase ($q > 0$)
- Disorder, **W**, decreased with T decrease ($q < 0$)
- As $T \rightarrow 0$, **W** → 1

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statements of the Second Law of Thermodynamics (roadmap #1)

1. Macroscopic properties of an isolated system eventually assume constant values (e.g. pressure in two bulbs of gas, becomes constant; two block of metal reach same T) →
2. It is impossible to construct a device that operates in cycles and that converts heat into work without producing some other change in the surroundings. *Kelvin's Statement [Raff p 157]; Carnot Cycle* →
3. It is impossible to have a natural process which produces no other effect than absorption of heat from a colder body and discharge of heat to a warmer body. *Clausius's Statement, refrigerator* →
4. In the neighborhood of any prescribed initial state there are states which cannot be reached by any adiabatic process
 ~ *Caratheodory's statement [Andrews p. 58]*


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goals of Carnot arithmetic (step 2 of roadmap)

1. Carnot cycle is "engine" that produces work from heat
2. Define efficiency:
 efficiency = (net work done by machine) / (heat energy input to machine)
3. Today, arithmetic manipulations of 1st Law results from ideal gas Carnot cycle (HW2 #10) to show that this efficiency depends only on the two temperatures at which heat is transferred to and from surroundings (the T_U of step 1 and T_L of step 3; the non-adiabatic paths)
4. Although for [reversible] Carnot cycle $\oint dq_{rev} \neq 0$ ^{WILL} but $\oint \frac{dq_{rev}}{T} = 0$ ^{WILL}

net work for cycle: $w_{net} > 0$
 engine DOES work on surroundings

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from Carnot cycle  <https://www.learnthermo.com/T1-tutorial/ch09/lesson-A/pg03.php>

for system in complete cycle:
 $\Delta U = 0$; $q > 0$; $w < 0$ (work DONE on surr) (**Prob #10e**)

$q > 0$ (q_{in}) at higher T_H ; $q < 0$ (q_{out}) at lower T_L

efficiency = $-w/q_{1 \rightarrow 2}$
 (how much total [net] work out (-sign) for heat in 1→2)

efficiency will depend on T_U and T_L

HW4 prob #22 ϵ is efficiency

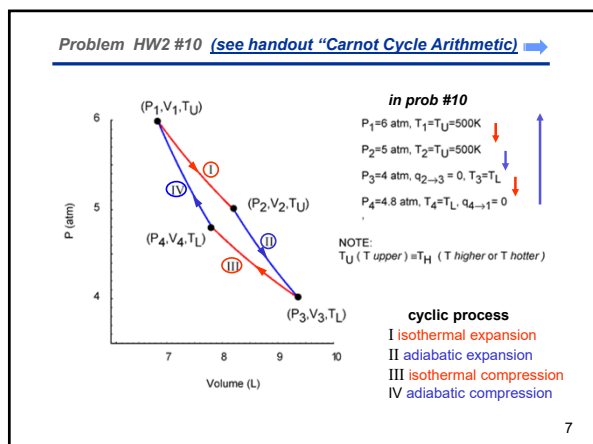
$$\epsilon = \frac{T_H - T_C}{T_H} \quad \text{or} \quad \epsilon = \frac{T_U - T_L}{T_U}$$

H=HOT C=COLD or U=UPPER L=LOWER

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let's go

- get $w_I + w_{II} + w_{III} + w_{IV} = w_{\text{total}}$
- get $q_I = q_{\text{input}}$

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Summary
 (see handout "Summary of Heat and Work for the Carnot Cycle Engines, Refrigerators, Heat Pumps")

general expressions for $(P_1, T_U) \xrightarrow{I} (P_2, T_U) \xrightarrow{II} (P_3, T_L) \xrightarrow{III} (P_4, T_L) \xrightarrow{IV} (P_1, T_U)$

ENGINE	q	w_{net}	w_{sur}	
I. isothermal expansion	$+nR T_U \ln \frac{P_1}{P_2} = 1.3$	$-nR T_U \ln \frac{P_1}{P_2} = 1.2$	$+nR T_U \ln \frac{P_1}{P_2}$	heat in at T_U work out
II. adiabatic expansion	0	$nC_v(T_U - T_L) = 2.4$	$-nC_v(T_U - T_L)$	work out
III. isothermal compression	$nR T_L \ln \frac{P_3}{P_4} = 3.36T.3$ $-nR T_L \ln \frac{P_3}{P_4}$	$-nR T_L \ln \frac{P_3}{P_4} = 3.26T.3$ $= nR T_L \ln \frac{P_1}{P_2}$	$-nR T_L \ln \frac{P_3}{P_4}$	heat lost at T_L work in
IV. adiabatic compression	0	$nC_v(T_U - T_L) = 4.4$	$-nC_v(T_U - T_L)$	work in
net gain/cost	$q_{\text{in}} = q_I$ $+nR T_U \ln \frac{P_1}{P_2}$		$w_{\text{sur}} = w_I + w_{II} + w_{III} + w_{IV}$ $nR(T_U - T_L) \ln \frac{P_1}{P_2}$	$\epsilon = w_{\text{sur}}/q_{\text{in}}$ $\epsilon = (T_U - T_L)/T_U$

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isothermal expansion at T_U (see handout "Carnot Cycle Arithmetic")

Step I: isothermal expansion, $T_U, V_1 \rightarrow V_2$

$$\Delta U_I = 0 \tag{1.1}$$

$$w_I = -nRT_U \ln \frac{V_2}{V_1} = nRT_U \ln \frac{P_2}{P_1} \tag{1.2}$$

$$q_I = -w_I = nRT_U \ln \frac{V_2}{V_1} = nRT_U \ln \frac{P_1}{P_2} \tag{1.3}$$

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step II: adiabatic reversible compression $T_U \rightarrow T_L$

$$V_2 = V_1 \left(\frac{T_U}{T_L} \right)^{\frac{1}{\gamma}} \quad V_3 = V_1 \left(\frac{P_1}{P_3} \right)^{\frac{1}{\gamma}} \tag{2.1}$$

$$T_2 = T_U \left(\frac{V_1}{V_2} \right)^{\gamma} \quad T_L = T_U \left(\frac{P_1}{P_2} \right)^{\frac{\gamma}{\gamma-1}} \tag{2.2}$$

$$P_2 = P_1 \left(\frac{V_1}{V_2} \right)^{\gamma} \quad P_2 = P_1 \left(\frac{T_U}{T_L} \right)^{\frac{\gamma}{\gamma-1}} \tag{2.3}$$

$$q_{II} = 0; \quad w_{II} = \Delta U \tag{2.4}$$

$$\Delta U_{II} = nC_v \Delta T = nC_v (T_L - T_U)$$

$$w_{II} = \Delta U_{II} = nC_v T_U \left(\left(\frac{V_2}{V_1} \right)^{\gamma} - 1 \right) = nC_v T_U \left(\left(\frac{P_1}{P_2} \right)^{\frac{\gamma}{\gamma-1}} - 1 \right) \tag{2.5}$$

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Step III: isothermal reversible compression at T_L

$$\Delta U_{III} = 0 \tag{3.1}$$

$$w_{III} = -nRT_L \ln \frac{V_4}{V_3} = nRT_L \ln \frac{P_3}{P_4} \tag{3.2}$$

$$q_{III} = -w_{III} = nRT_L \ln \frac{V_4}{V_3} = nRT_L \ln \frac{P_3}{P_4} \tag{3.3}$$

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Step IV: adiabatic reversible compression ($T_L \rightarrow T_U$)

$$V_i = V_i \left(\frac{T_U}{T_L} \right)^{\frac{1}{\gamma}} \quad V_f = V_i \left(\frac{P_i}{P_f} \right)^{\frac{1}{\gamma}} \quad (4.1)$$

$$T_i = T_i \left(\frac{V_i}{V_f} \right)^{\gamma} \quad T_f = T_f \left(\frac{P_i}{P_f} \right)^{\frac{\gamma}{1-\gamma}} \quad (4.2)$$

$$P_i = P_i \left(\frac{V_i}{V_f} \right)^{\gamma} \quad P_f = P_f \left(\frac{T_i}{T_f} \right)^{\frac{\gamma}{1-\gamma}} \quad (4.3)$$

P.V.T relationships for adiabatic reversible process

$$q_{IV} = 0; \quad w_{IV} = \Delta U_{IV} \quad (4.4)$$

$$\Delta U_{IV} = n\bar{C}_V \Delta T = n\bar{C}_V (T_U - T_L)$$

$$w_{IV} = \Delta U_{IV} = n\bar{C}_V T_U \left[1 - \left(\frac{V_i}{V_f} \right)^{\gamma} \right] = n\bar{C}_V T_U \left[1 - \left(\frac{P_i}{P_f} \right)^{\frac{1}{\gamma}} \right] \quad (4.5)$$

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note: $w_{II} = -w_{IV}$ the two adiabatic steps have opposite work out \leftrightarrow work in

$$\Delta U_{II} = n\bar{C}_V (T_L - T_U) = w_{II} \quad \longleftrightarrow \quad w_{IV} = n\bar{C}_V (T_U - T_L) = \Delta U_{IV}$$

$$w_{II} = n\bar{C}_V T_U \left[\left(\frac{P_3}{P_2} \right)^{\frac{\gamma}{1-\gamma}} - 1 \right] = n\bar{C}_V T_U \left[\left(\frac{4}{5} \right)^{\frac{\gamma}{1-\gamma}} - 1 \right] = n\bar{C}_V T_U \left[(0.8)^{\frac{\gamma}{1-\gamma}} - 1 \right]$$

$$w_{IV} = n\bar{C}_V T_U \left[1 - \left(\frac{P_1}{P_2} \right)^{\frac{\gamma}{1-\gamma}} \right] = n\bar{C}_V T_U \left[1 - \left(\frac{4.8}{6} \right)^{\frac{\gamma}{1-\gamma}} \right] = n\bar{C}_V T_U \left[1 - (0.8)^{\frac{\gamma}{1-\gamma}} \right]$$

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and now for the TOTAL cycle (T_U and T_L ; and P_1 and P_2 given)

$$w_{total} = w_I + w_{II} + w_{III} + w_{IV}$$

$$w_{II} = -w_{IV} \implies w_{total} = w_I + w_{III}$$

$$w_I = nRT_U \ln \frac{P_2}{P_1}$$

$$w_{III} = nRT_L \ln \frac{P_1}{P_2} \quad \text{with} \quad P_3 = P_2 \left(\frac{T_U}{T_L} \right)^{\frac{1}{\gamma}} \quad \text{and} \quad P_4 = P_1 \left(\frac{T_U}{T_L} \right)^{\frac{1}{\gamma}}$$

$$w_{III} = nRT_L \ln \left(\frac{P_1 \left(\frac{T_U}{T_L} \right)^{\frac{1}{\gamma}}}{P_2 \left(\frac{T_U}{T_L} \right)^{\frac{1}{\gamma}}} \right) = nRT_L \ln \left(\frac{P_1}{P_2} \right) = -nRT_L \ln \left(\frac{P_2}{P_1} \right)$$

$$w_{total} = nR(T_U - T_L) \ln \frac{P_2}{P_1}$$

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and NOW EFFICIENCY ϵ

efficiency: $\epsilon = \frac{\text{total work done ON SURROUNDINGS}}{\text{heat INPUT}}$

$$\epsilon = \frac{-w_{total}}{q_I}$$

$$w_{total} = nR(T_U - T_L) \ln \frac{P_2}{P_1}$$

$$q_I = -w_I = nRT_U \ln \frac{P_1}{P_2} = -nRT_U \ln \frac{P_2}{P_1}$$

and

$$\epsilon = \frac{-nR(T_U - T_L) \ln \frac{P_2}{P_1}}{-nRT_U \ln \frac{P_2}{P_1}} = \frac{(T_U - T_L)}{T_U} = 1 - \frac{T_L}{T_U}$$

only $q_I \equiv q_{UPPER}$
 q_{III} is wasted heat lost to surroundings at T_L as thermal pollution

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Summary
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general expressions for $(P_1, T_U) \xrightarrow{I} (P_2, T_U) \xrightarrow{II} (P_3, T_L) \xrightarrow{III} (P_4, T_L) \xrightarrow{IV} (P_1, T_U)$

ENGINE	q	w_{sys}	w_{sur}	
I. isothermal expansion	$+nRT_U \ln \frac{P_1}{P_2}$ 1.3	$-nRT_U \ln \frac{P_1}{P_2}$ 1.2	$+nRT_U \ln \frac{P_1}{P_2}$	heat in at T_U work out
II adiabatic expansion	0	$n\bar{C}_V (T_U - T_L)$ 2.4	$-n\bar{C}_V (T_U - T_L)$	work out
III. isothermal compression	$nRT_L \ln \frac{P_1}{P_2} = -nRT_L \ln \frac{P_2}{P_1}$ 3.3&3.4	$-nRT_L \ln \frac{P_1}{P_2} = nRT_L \ln \frac{P_2}{P_1}$ 3.2&3.3	$-nRT_L \ln \frac{P_1}{P_2}$	heat lost at T_L work in
IV. adiabatic compression	0	$n\bar{C}_V (T_U - T_L)$ 4.4	$-n\bar{C}_V (T_U - T_L)$	work in
net gain/cost	$q_{in} = q_I$ $+nRT_U \ln \frac{P_1}{P_2}$		$w_{total} = w_I + w_{II} + w_{III} + w_{IV} =$ $nR(T_U - T_L) \ln \frac{P_2}{P_1}$	$\epsilon = w_{total}/q_{in}$ $\epsilon = (T_U - T_L)/T_U$

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some limits on efficiency of ideal engine

$$\epsilon = \frac{(T_U - T_L)}{T_U} = 1 - \frac{T_L}{T_U}$$

lim $\epsilon = 0$
 $T_L \rightarrow T_U$ must have q in at higher and q out at lower T

lim $\epsilon = 1$
 $T_U \rightarrow \infty$ perfect efficiency at finite temperatures only for $T_{LOWER} = 0$ K

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roadmap for second law

- ✓ 1. Phenomenological statements (what is ALWAYS observed)
- ✓ 2. Ideal gas Carnot [reversible] cycle efficiency of heat → work (Carnot cycle transfers heat only at T_U and T_L)
3. Any cyclic engine operating between T_U and T_L must have an equal or lower efficiency than Carnot OR VIOLATE one of the phenomenological statements (observations)
4. Generalize Carnot to any reversible cycle (E&R fig 5.4)
5. Show that for this REVERSIBLE cycle



$$q_U + q_L \neq 0 \quad (\text{dq inexact differential } \oint dq \neq 0)$$
 but

$$\frac{q_U}{T_U} + \frac{q_L}{T_L} = 0 \quad (\text{something VERY VERY special about } \frac{dq_{rev}}{T}; \oint \frac{dq_{rev}}{T} = !!!)$$
6. S, entropy and spontaneous changes

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End of Lecture 9

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take home messages (lecture 8) and statements of 2nd Law (Andrews)

Two bodies 500 molecules each
 one at $T_A = .63 \times T_f$
 the other at $T_B = 1.37 \times T_f$

.63 T_f

1.37 T_f

bring into thermal contact

STAY COLD-HOT ??

or

EQUILIBRIUM ??

$T_E = T_f$

↑

DISORDER

↑

~ microstates

$\approx 10^{556}$

$= 10^{205} \times 10^{354}$

(relative)

Q T F GT

↑

DISORDER

↑

↑

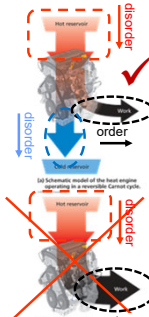
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E&R_{4th} fig 5.14 p. 127 [fig 5.3, p. 89]_{E&R3rd}

These considerations on the efficiency of reversible heat engines led to the Kelvin-Planck formulation of the second law of thermodynamics:

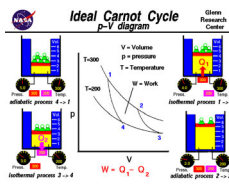
It is impossible for a system to undergo a cyclic process whose sole effects are the flow of heat into the system from a heat reservoir and the performance of an equal amount of work by the system on the surroundings.

Figure 5.14 Carnot heat engine cycle. (a) A schematic model of a heat engine operating in a reversible Carnot cycle. (b) The second law of thermodynamics asserts that it is impossible to construct a heat engine that operates using a single heat reservoir and converts the heat withdrawn from the reservoir into work with 100% efficiency as shown.



Ideal Carnot Cycle

p-V diagram



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take home messages (lecture 8) and statements of 2nd Law (Clausius)

Two bodies 500 molecules each
 one at $T_A = .63 \times T_f$
 the other at $T_B = 1.37 \times T_f$

.63 T_f

1.37 T_f

(keep) transferring heat from colder body to warmer body

$T_E = T_f$

↑

DISORDER

↑

~ microstates

$\approx 10^{556}$

$10^{205} \times 10^{354}$

(relative)

Q T F GT

↑

DISORDER

↑

↑

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