Chemistry 163B Winter 2020

Lectures 2-3

Heat and Work

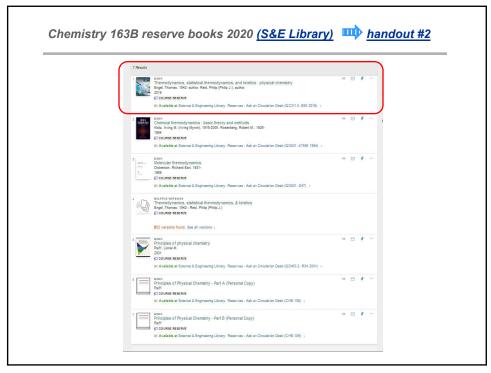
1

Menu: for TODAY(s)



- Heat (q), Work (w) and the First Law of Thermodynamics
 ∆U≡q+w
- Example calculations of q, w, U for reversible and irreversible expansions of ideal gas (comparisons and 'take home messages')
- · Molecular picture of heat and work energy transfers

2



3

a few definitions ,etc. to get us going



- df is 'differential' representing an infinitesimal change
- to calculate Δf a finite change $\Delta f = \int_{initial}^{final} df$



 if f is a state function, the integral around a closed (cyclic) path (finish=start) is

$$\oint df = 0$$



4

a few "FACTS OF (thermodynamic) LIFE"



if the 'system' goes from one state to another by differing paths

initial state — final state



· the change in some aspect of the system can either be

independent of the path taken
 (depends only on the initial and final values of the state variables)

depend on the path taken

if change is path independent ⇒ exact differential,

there is underlying **STATE FUNCTION**

if change is path dependent ⇒ inexact differential,
 no underlying property of the state of system

more in Lecture #5 (math !!)

5

5

all transfers of energy can be classified as heat (q) or work (w)

We will see (and demonstrate):



- heat (q) and work (w) are NOT PROPERTIES of a system but correspond to TRANSFERS of energy into (+)_{sys} or out of (-)_{sys} the system [e.g. there is no underlying property "heat" of a system"]
- when a system goes from an initial to a final state,
 e.g. (P_i, V_i, T_i) ⇒ (P_f, V_f, T_f) the values of **q** and **w** will **DEPEND ON THE PATH** taken between the states
- a small changes in a **path-dependent** quantities are **INEXACT DIFFERENTIALS** indicated by e.g. (t)q and (t)w

6

the FIRST LAW of thermodynamics



- U is symbol for INTERNAL ENERGY of a system

•
$$\Delta U = q + w + \chi$$
; $dU = dq + dw + \sum_{i=1}^{n} dn_i$

dn_i=0 closed system,

- $\Delta U_{sys} = -\Delta U_{surr}$ (conservation of energy)
- U is state function
- ∆U does not depend on path (only on initial and final states)

equivalent statements

dU is an exact differential

7

heat capacity (E&R section 2.11)[4th]

 $\frac{d^{2}q}{dT} = C \qquad heat \ capacity \left[J K^{-1} \right] \ extensive$

the amount (transfer) of heat required to raise substance 1K

 $\frac{d\overline{q}}{dT} = n\overline{C} \quad molar \ heat \ capacity \ [J \ mol^{-1} \ K^{-1}] \ intensive$

the amount (transfer) of heat requires to raise 1 mol substance 1K

 \overline{C} generally depends on T and conditions for example ideal monatomic gas

molar heat capacity at constant volume \bar{C}_{V} : molar heat capacity at constant pressure $\overline{C}_p = \frac{3}{2}$

 $E = \frac{3}{2} nRT (true!!)$ but why does it take more heat to raise T at constant P than at constant V ??

transfers of energy: HEAT and work (sec 2.2)[4th]

change of energy by heat transfer

$$\vec{d} q = C dT = n \overline{C} dT$$
 (C is extensive,
 $q = \int_{path} \vec{d} q = \int_{path} n \overline{C} dT$ \overline{C} is intensive)

C will generally depend on T and path (i.e. conditions)

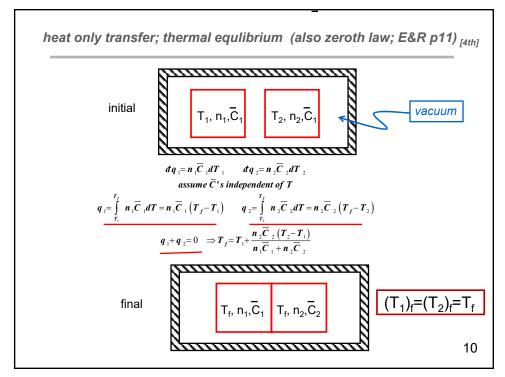
 $q > 0 \Rightarrow$ energy (heat) gained by system (endothermic) cools surroundings



 $q < 0 \Rightarrow$ energy (heat) lost by system (exothermic) warms surroundings

9

9



change of energy by work done ON system

energy into system

derive in a moment.

$$d w = d w_{PV} + d w_{other}$$

$$d w_{PV} = -P_{\underline{ext}} dV$$

$$w_{PV} = \int d w_{PV} = \int -P_{\underline{ext}} dV$$
e.g. electrical, gravitational, etc.

 $w > 0 \implies$ energy gained by system (work done ON system) $w < 0 \implies$ energy lost by system

(work done ON surroundings)

11

11

processes: definitions of constraints

isolated

q=0; w=0

isothermal

∆T=0

adiabatic

q=0

• "against constant

pressure"

 P_{ext} = const

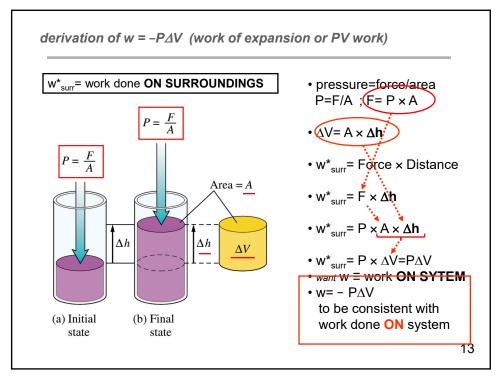
• reversible process $P_{int} = P_{ext}$

a (ideal) process that proceeds so slowly that an infinitesimal change of conditions causes the process to proceed in the opposite (reverse) direction

• irreversible process

all other (real) processes proceeding at finite rate

12



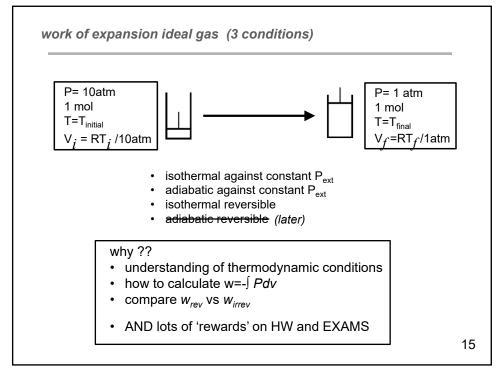
13

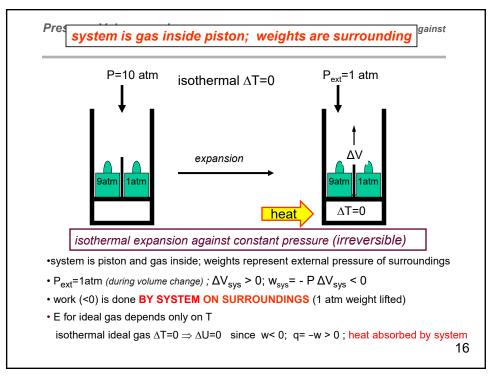
ideal gas and energy, heat, work

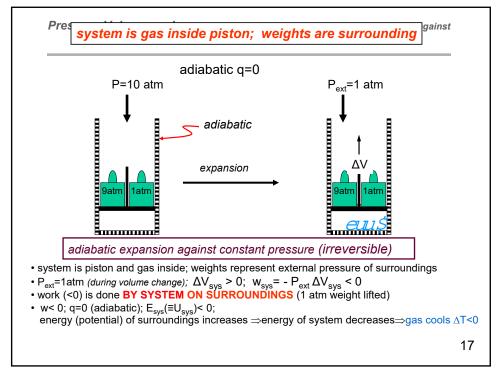
for IDEAL GAS

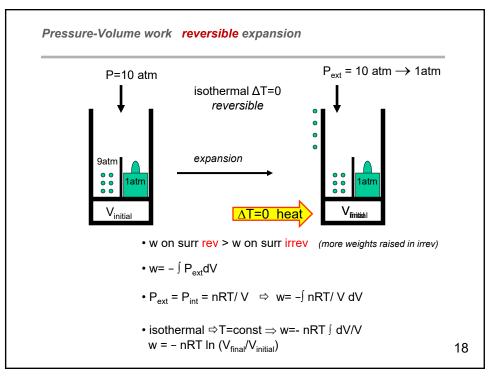
- U(≡E) depends ONLY on T (ideal gas, previous class)
- isothermal, ΔT=0,
 - ∆U=0=q+w (ideal gas)
 - q=-w
- adiabatic: q=0, ΔU=w (in general)
- monatomic ideal gas
 - U = (3/2) n RT
 - C_V = (3/2) n R (~ prove later)
 - $C_P = (5/2) \text{ n R (prove later)}$

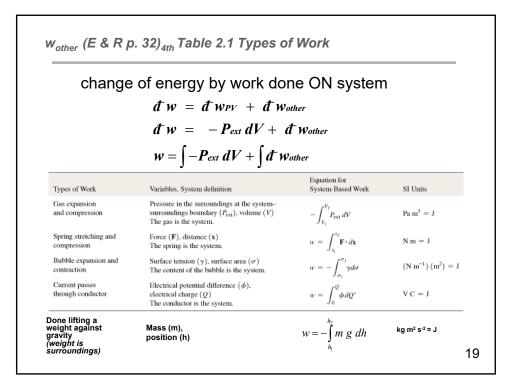
14

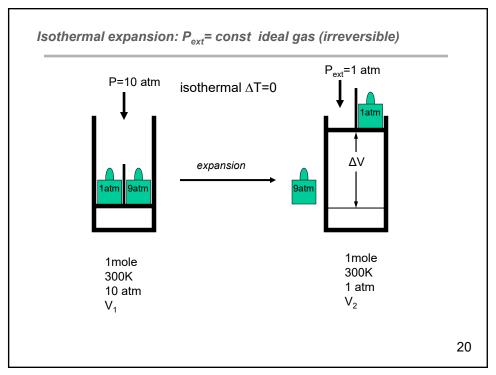


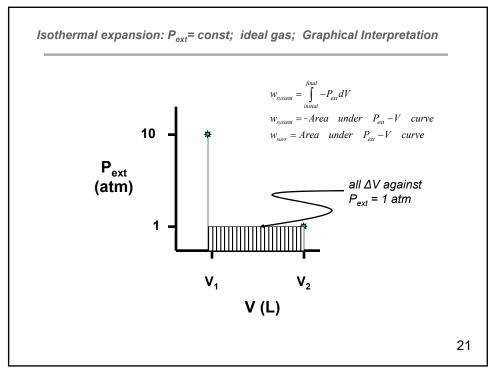


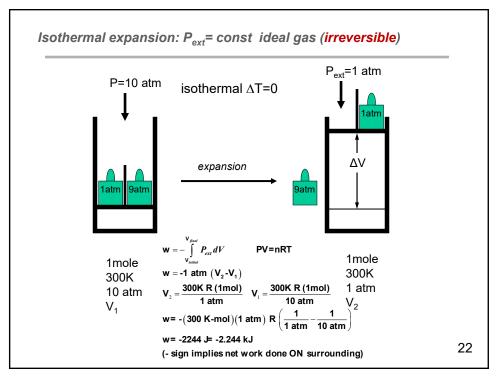


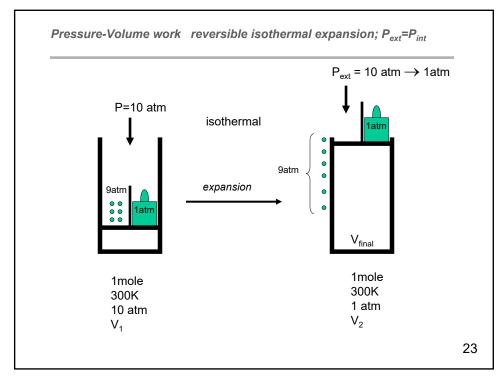


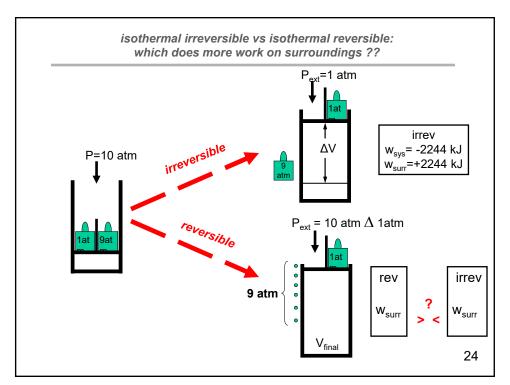


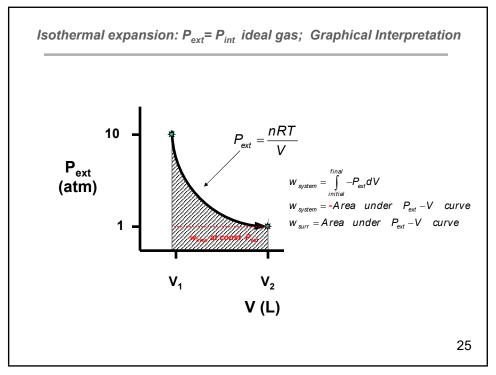


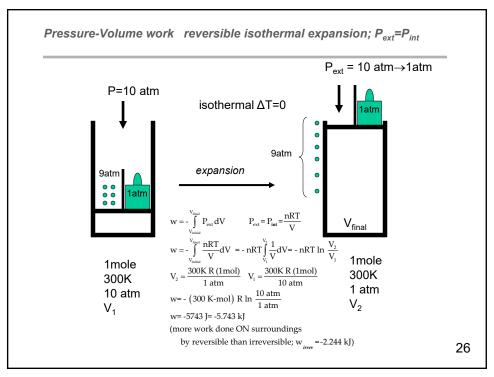


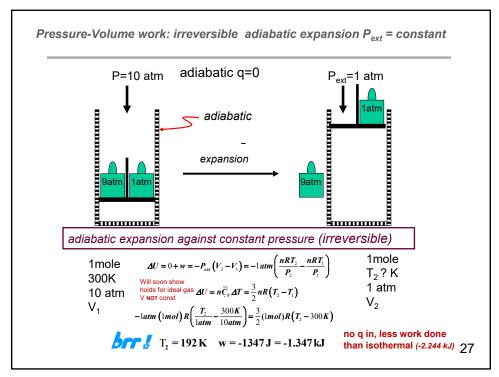












27



· total energy:

sum of [number of particles in energy level (v) × energy of that level (ϵ)] $E=\Sigma_n v_n \varepsilon_n$

• heat only (no work):

redistributes particles among levels

 $\text{q>}0 \Rightarrow \text{more particles to higher energy levels}$

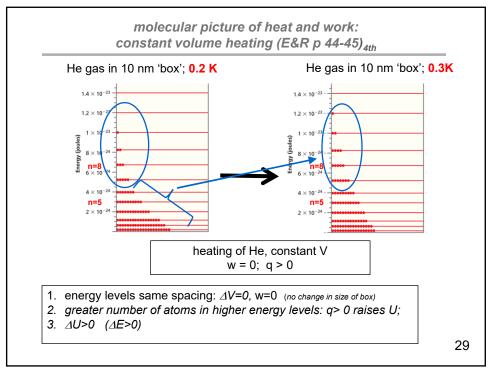
q<0 ⇒ more particles to lower energy levels

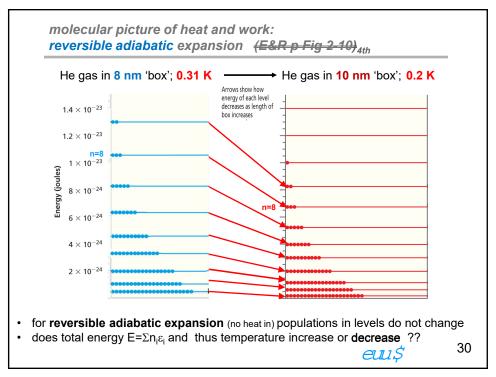
P-V work only (adiabatic)

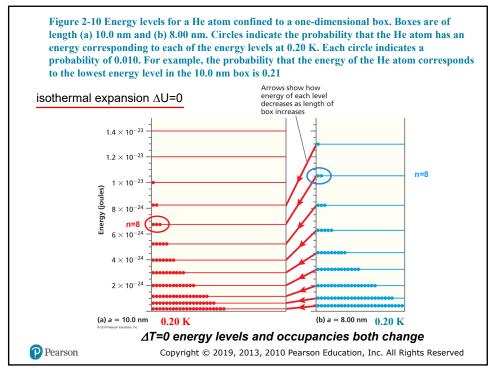
particle in 3D cube: CHEM 163A DON'T FRET $\varepsilon_n \approx \frac{3h^2}{2}n^2$ where n is an integer (1, 2, ...)

larger volume $(w<0) \Rightarrow$ energy levels lower and get more closely spaced smaller volume $(w>0) \Rightarrow$ energy levels higher and less closely spaced

28







31

Menu: for TODAY(s)



- ✓ Heat (q), Work (w) and the First Law of Thermodynamics
 △U≡q+w
- ✓ Example calculations of q, w, U for reversible and irreversible expansions of ideal gas (comparisons and 'take home messages')
- ✓ Molecular picture of heat and work energy transfers

32

End of Lectures 2-3