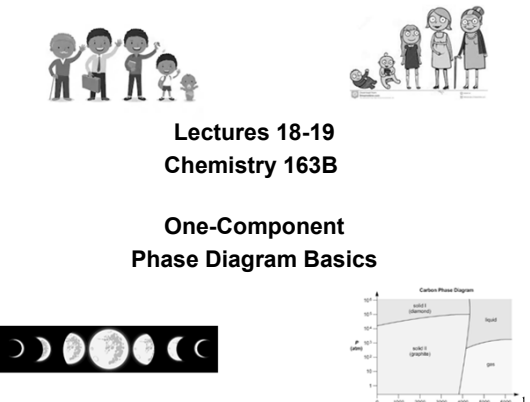


Chemistry 163B, Winter 2020
Lectures 18-19 Introduction to Phase Diagrams



Lectures 18-19
Chemistry 163B

One-Component
Phase Diagram Basics

qualitative factors in phase changes

melting
solid \rightleftharpoons liquid

freezing/fusion

vaporization
liquid \rightleftharpoons gas **ENDOTHERMIC**

condensation **EXOTHERMIC**

sublimation
solid \rightleftharpoons gas

deposition

2

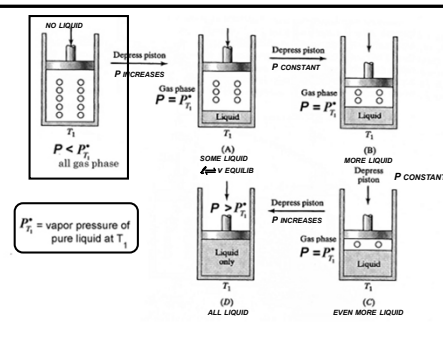
vapor pressure over PURE liquid (notation)

$P^\bullet \equiv P^* \equiv P^0$

Gene (many others) E&R Raff

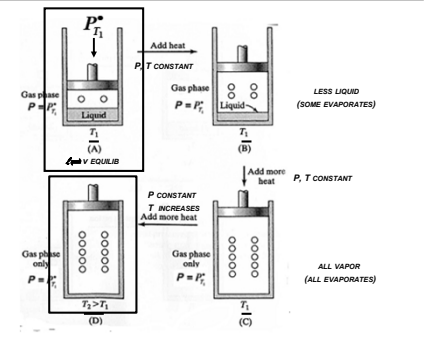
3

gas \rightleftharpoons liquid as pressure increases (vary P, const T)



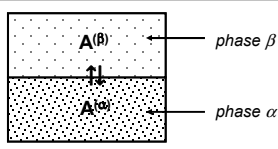
4

liquid \rightleftharpoons vapor as heat added (vary T, const P)



5

dG for phase change at constant T,P X^(y) property X, phase y



same T,P for each phase

$$dG = -SdT + VdP + \sum_i \mu_i dn_i$$

one component 'A' in phases α and β constant T,P

$$dG_{T,P} = \mu_A^{(\alpha)} dn_A^{(\alpha)} + \mu_A^{(\beta)} dn_A^{(\beta)}$$

$$dn_A^{(\beta)} = -dn_A^{(\alpha)}$$

$$dG_{T,P} = (\mu_A^{(\alpha)} - \mu_A^{(\beta)}) dn_A^{(\alpha)}$$

6

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at equilibrium $\mu^{(\alpha)} = \mu^{(\beta)}$; μ is ESCAPING TENDENCY

$$dG_{T,P} = (\mu_A^{(\alpha)} - \mu_A^{(\beta)}) dn_A^{(\alpha)}$$

at equilibrium $dG_{T,P} = 0$

$$\mu_A^{(\alpha)} = \mu_A^{(\beta)}$$

for spontaneity $dG_{T,P} < 0$

$$dG_{T,P} = (\mu_A^{(\alpha)} - \mu_A^{(\beta)}) dn_A^{(\alpha)} < 0$$

$\mu_A^{(\alpha)} > \mu_A^{(\beta)} \Rightarrow dn_A^{(\alpha)} < 0$ **molecules lost from phase α**

$\mu_A^{(\beta)} > \mu_A^{(\alpha)} \Rightarrow dn_A^{(\alpha)} > 0$ **molecules gained by phase α**

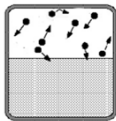
$\mu_A^{(\alpha)}$ is the **ESCAPING TENDENCY** for molecules in phase α

high $\mu \rightarrow$ low μ
hyper \rightarrow mellow

7

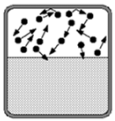
Question:

can my pressure cooker heat water to 200C without exploding?



T=298 K
P*=0.032 bar

➡



T=473 K
P*=15.5 bar

https://www.engineeringtoolbox.com/water-vapor-saturation-pressure-d_599.html

https://www.engineeringtoolbox.com/water-properties-d_1573.html

8

phase equilibrium one-component systems (i.e pure substances)

A(α) \rightleftharpoons A(β)

at equilibrium $\Delta\mu = 0 \Rightarrow \mu_A^{(\alpha)} = \mu_A^{(\beta)}$

How can P and T covary to maintain equilibrium?

$\mu^{(\alpha)}(T_1, P_1) = \mu^{(\beta)}(T_1, P_1)$

T and P covary

$d\mu^{(\alpha)}$
↓
 $d\mu^{(\beta)}$

T and P covary

$\mu^{(\alpha)}(T_1, P_1) = \mu^{(\beta)}(T_1, P_1)$

before

➡

$\mu^{(\alpha)}(T_2, P_2) = \mu^{(\beta)}(T_2, P_2)$

after

9

conditions for remaining at phase equilibrium (one-component), covary T and P

$$\mu^{(\alpha)}(T_1, P_1) = \mu^{(\beta)}(T_1, P_1) \xrightarrow{dT, dP} \mu^{(\alpha)}(T_2, P_2) = \mu^{(\beta)}(T_2, P_2)$$

$$d\mu^{(\alpha)} = -\bar{S}^{(\alpha)} dT^{(\alpha)} + \bar{V}^{(\alpha)} dP^{(\alpha)} = -\bar{S}^{(\beta)} dT^{(\beta)} + \bar{V}^{(\beta)} dP^{(\beta)} = d\mu^{(\beta)}$$

with

$$T^{(\alpha)} = T^{(\beta)} = T \quad dT^{(\alpha)} = dT^{(\beta)} = dT$$

$$P^{(\alpha)} = P^{(\beta)} = P \quad dP^{(\alpha)} = dP^{(\beta)} = dP$$

$$-\bar{S}^{(\alpha)} dT + \bar{V}^{(\alpha)} dP = -\bar{S}^{(\beta)} dT + \bar{V}^{(\beta)} dP$$

$$(\bar{S}^{(\beta)} - \bar{S}^{(\alpha)}) dT = (\bar{V}^{(\beta)} - \bar{V}^{(\alpha)}) dP$$

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phase equilibrium (one-component)

$$(\bar{S}^{(\beta)} - \bar{S}^{(\alpha)}) dT = (\bar{V}^{(\beta)} - \bar{V}^{(\alpha)}) dP$$

$\left(\frac{dP}{dT}\right)_{\text{phase equilibrium}} = \frac{\Delta\bar{S}_\phi}{T\Delta\bar{V}_\phi}$

← $\phi = \text{phase change}$

eqn. 8.13 E&R_{4th}

since phase change is an equilibrium (reversible) process

$$\Delta\bar{S}_\phi = \frac{\Delta\bar{H}_\phi}{T}$$

$$\left(\frac{dP}{dT}\right)_{\text{phase equilibrium}} = \frac{\Delta\bar{H}_\phi}{T\Delta\bar{V}_\phi}$$

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1. application to liquid \rightleftharpoons gas (vapor) or solid \rightleftharpoons gas

$$\left(\frac{dP}{dT}\right)_{\text{phase equilibrium}} = \frac{\Delta\bar{H}_\phi}{T\Delta\bar{V}_\phi}$$

vaporization – condensation liquid \rightleftharpoons gas (vapor)

or

sublimation – deposition solid \rightleftharpoons gas

\bar{V}_{solid} and \bar{V}_{liquid} are small compared to \bar{V}_{vapor} $\Delta\bar{V}_\phi \approx \bar{V}_{\text{vapor}}$

assume ideal gas $\Delta\bar{V}_\phi = \bar{V}_{\text{vapor}} = \frac{RT}{P}$

$$\left(\frac{dP}{dT}\right)_{\text{or } \rightleftharpoons \text{g}} = \frac{\Delta\bar{H}_\phi}{T\left(\frac{RT}{P}\right)} = \frac{P\Delta\bar{H}_\phi}{RT^2}$$

$$\left(\frac{d(\ln P)}{dT}\right) = \frac{\Delta\bar{H}_{\text{vap or sub}}}{RT^2}$$

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Lectures 18-19 Introduction to Phase Diagrams

I. application to liquid ⇌ gas (vapor) or solid ⇌ gas

$$\left(\frac{dP}{dT}\right)_{\text{equilibrium}} = \frac{\Delta\bar{H}_\phi}{T} = \frac{P\Delta\bar{H}_\phi}{RT^2}$$


$$\left(\frac{d(\ln P)}{dT}\right)_{\text{equilibrium}} = \frac{\Delta\bar{H}_{\text{vap or sub}}}{RT^2}$$

Clausius-Clapeyron
≈ eqn. 8.19 E&R_{4th}

So just like before:

for s or l ⇌ gas $K_p = (a_{\text{gas}}/a_{s,l}) \approx (P_{\text{gas}}/1)$

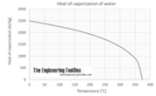
$$\left(\frac{d(\ln K_p)}{dT}\right)_{\text{equilibrium}} = \frac{\Delta\bar{H}_{\text{vap or sub}}}{RT^2}$$


 midterm 2

13

I. application to liquid ⇌ gas (vapor) or solid ⇌ gas

$$\left(\frac{d(\ln P)}{dT}\right)_{\text{equilibrium}} = \frac{\Delta\bar{H}_{\text{vap or sub}}}{RT^2}$$



$$\int_{P_1}^{P_2} d(\ln P) = \int_{T_1}^{T_2} \frac{\Delta\bar{H}_\phi}{RT^2} dT \quad (\text{assume } \Delta\bar{H} \text{ independent of } T \text{ sometimes??})$$

$$\ln\left(\frac{P_2}{P_1}\right) = -\frac{\Delta\bar{H}_\phi}{R} \left[\frac{1}{T_2} - \frac{1}{T_1}\right]$$

E&R_{4th} eqn 8.20 where ϕ is vaporization similar for sublimation

application to problems: normal b.p. (1 atm), standard b.p. (1 bar)

to get vapor pressure given T_{boiling}° and ΔH_{vap} :

at $T_1 = T_{\text{bp}}^\circ$ $P_1 = P_{\text{vap}} = 1 \text{ atm}$

$$\ln\left(\frac{P_{\text{vap}}(T)}{1 \text{ atm}}\right) = -\frac{\Delta\bar{H}_{\text{vap}}}{R} \left[\frac{1}{T} - \frac{1}{T_{\text{bp}}}\right] = \frac{\Delta\bar{H}_{\text{vap}}}{R} \left[\frac{1}{T_{\text{bp}}} - \frac{1}{T}\right]$$

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I. application to liquid ⇌ gas (vapor) or solid ⇌ gas

application to problems: to get T_{boiling} when $P_{\text{atm}} \neq 1 \text{ atm}$:

$$\ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right) = -\frac{\Delta\bar{H}_{\text{vap}}}{R} \left[\frac{1}{T_{\text{bp}}} - \frac{1}{T_{\text{bp}}^\circ}\right]$$

$$T_{\text{bp}}^\circ \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right) = -\frac{\Delta\bar{H}_{\text{vap}}}{R} \left[\frac{T_{\text{bp}}^\circ}{T_{\text{bp}}} - 1\right]$$

$$T_{\text{bp}}^\circ \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right) - \frac{\Delta\bar{H}_{\text{vap}}}{R} = -\frac{\Delta\bar{H}_{\text{vap}}}{R} \left[\frac{T_{\text{bp}}^\circ}{T_{\text{bp}}}\right]$$

$$\left[T_{\text{bp}}^\circ \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right) - \frac{\Delta\bar{H}_{\text{vap}}}{R}\right] \left[\frac{R}{\Delta\bar{H}_{\text{vap}}}\right] = -\left[\frac{T_{\text{bp}}^\circ}{T_{\text{bp}}}\right]$$

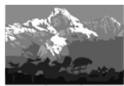
$$\left[\frac{T_{\text{bp}}}{T_{\text{bp}}^\circ}\right] = \frac{1}{1 - \frac{RT_{\text{bp}}^\circ}{\Delta\bar{H}_{\text{vap}}} \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right)}$$

15

I. application to liquid ⇌ gas (vapor) or solid ⇌ gas


application to problems: to get T_{boiling} when $P_{\text{atm}} \neq 1 \text{ atm}$:

$$\left[\frac{T_{\text{bp}}}{T_{\text{bp}}^\circ}\right] = \frac{1}{1 - \frac{RT_{\text{bp}}^\circ}{\Delta\bar{H}_{\text{vap}}} \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right)}$$



Denver: elev=1610 m $P=0.822 \text{ atm}$

$P < 1 \text{ atm}$ $\ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right) < 0 \Rightarrow \left[\frac{T_{\text{bp}}}{T_{\text{bp}}^\circ}\right] < 1 \Rightarrow T_{\text{bp}} < T_{\text{bp}}^\circ$



Death Valley: elev = -82.5 m, $P=1.010 \text{ atm}$

$P > 1 \text{ atm}$ $\ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right) > 0 \Rightarrow \left[\frac{T_{\text{bp}}}{T_{\text{bp}}^\circ}\right] > 1 \Rightarrow T_{\text{bp}} > T_{\text{bp}}^\circ$

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Denver vs Death Valley

$$\left[\frac{T_{\text{bp}}}{T_{\text{bp}}^\circ}\right] = \frac{1}{1 - \frac{RT_{\text{bp}}^\circ}{\Delta\bar{H}_{\text{vap}}} \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right)}$$

$$\left[\frac{T_{\text{bp}}}{373}\right] = \frac{1}{1 - \frac{(8.3245 \text{ J K}^{-1} \text{ mol}^{-1})(373 \text{ K})}{40.65 \times 10^3 \text{ J mol}^{-1}} \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right)} = \frac{1}{1 - 0.0763 \ln\left(\frac{P_{\text{atm}}}{1 \text{ atm}}\right)}$$

Denver: elev 1610 m $P=0.822 \text{ atm}$ $T_{\text{bp}}=367.5 \text{ K}$
 Sea level: elev 0 m $P=1.00 \text{ atm}$ $T_{\text{bp}}=373 \text{ K}$
 Death Valley: elev -82.5 m $P=1.01 \text{ atm}$ $T_{\text{bp}}=373.28 \text{ K}$

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II. application to solid ⇌ liquid

$$\left(\frac{dP}{dT}\right)_{\text{phase equilibrium}} = \frac{\Delta\bar{H}_\phi}{T\Delta\bar{V}_\phi} \quad T_{\text{melt}} \text{ for phase equilibrium at } P = 1 \text{ atm} \quad \text{what is } T_{\text{melt}} \text{ at other pressures?}$$

$$\frac{dT}{T} = \frac{\Delta\bar{V}_\phi}{\Delta\bar{H}_{\text{melting}}} dP \Rightarrow \ln\left(\frac{T_{\text{melting}}}{T_{\text{melting}}^\circ}\right) = \frac{\Delta\bar{V}_\phi}{\Delta\bar{H}_{\text{melting}}} [P - 1 \text{ atm}]$$

$$\ln\left(\frac{T_{\text{melting}}}{T_{\text{melting}}^\circ}\right) = \frac{\bar{V}_{\text{liquid}} - \bar{V}_{\text{solid}}}{\Delta\bar{H}_{\text{melting}}} [P - 1 \text{ atm}]$$

will increased pressure raise or lower T_{melting} ?

$\Delta H_{\text{melting}} > 0$

(usual) $V_{\text{liquid}} > V_{\text{solid}}$ T_{melting} increases

(when??) $V_{\text{liquid}} < V_{\text{solid}}$ T_{melting} decreases

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Lectures 18-19 Introduction to Phase Diagrams

phase rule one-component system (save proof for later; $f=c-p+2$, $c=1$)

f = degrees of freedom
 p = phases simultaneously present

2 variables : T, P (same for each phase)

$p-1$ restrictions: $\mu^{(\alpha)} = \mu^{(\beta)} = \mu^{(\gamma)} = \dots$

f : degrees of freedom = (variables-restrictions)
 $f = 2 - (p-1) = 3-p$

$f = 3-p$

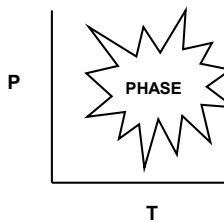
1 phase: T,P vary independently

2 phases present: T and P covary

3 phases present: fixed T and P

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phase diagrams one component: phase vs (P, T)



BE[A]WARE: when we study multicomponent phase diagrams the axis variables may not be P,T

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HW7 #45 (phase diagram but note log P axis scaling)

45. E8Fuu PB.1
 Within what range can you restrict the values of P and T if the following information is known about CO₂? Use Figure 8.12 to answer this question.
 Note: The critical point is at 31°C, 73 atm, and P_{tr} is 5.1 atm.

- As the temperature is increased, the solid is first converted to the liquid and subsequently to the gaseous state.
- As pressure on a cylinder containing pure CO₂ is increased from 8 to 80 atm, no interface separating liquid and gaseous phases is observed, (over the 8 atm range affect from E8Fuu CO2 TRIP VALUE, if makes more sense)
- Solid, liquid, and gas phases coexist at equilibrium.
- An increase in pressure from 10 to 50 atm converts the liquid to the solid.
- An increase in temperature from -80 to 20 °C converts a solid to a gas with no intermediate liquid phase.

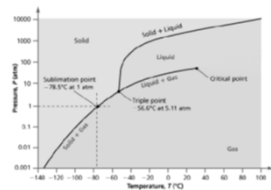
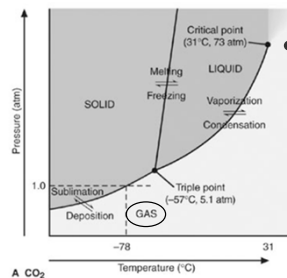


Figure 8.12 The P-T phase diagram for CO₂

21

phase diagrams ($f=3-p$)

“state” or “phase” as a function of P, T

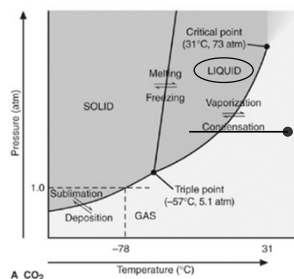


high T
 gas
 1 phase, $f=2$
 vary both T,P

22

phase diagrams ($f=3-p$)

“state” or “phase” as a function of P, T

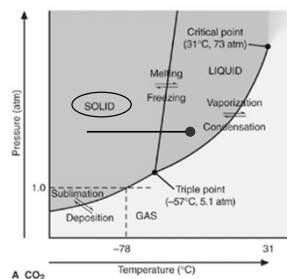


lower T
 liquid
 1 phase, $f=2$
 vary both T,P

23

phase diagrams ($f=3-p$)

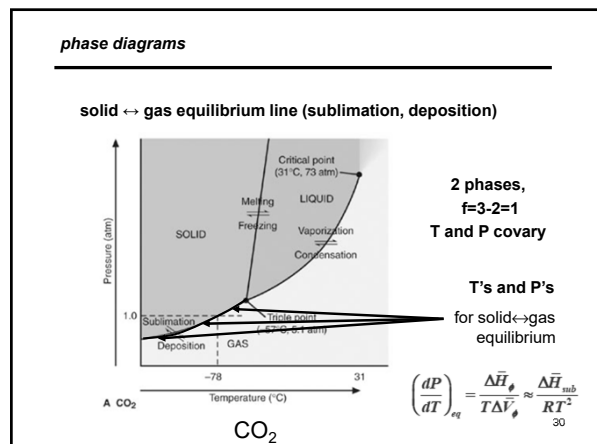
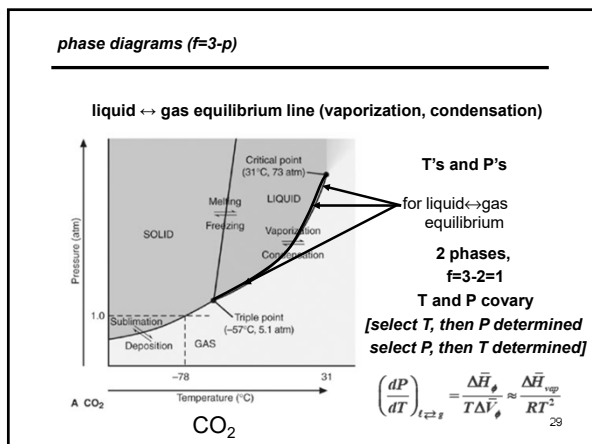
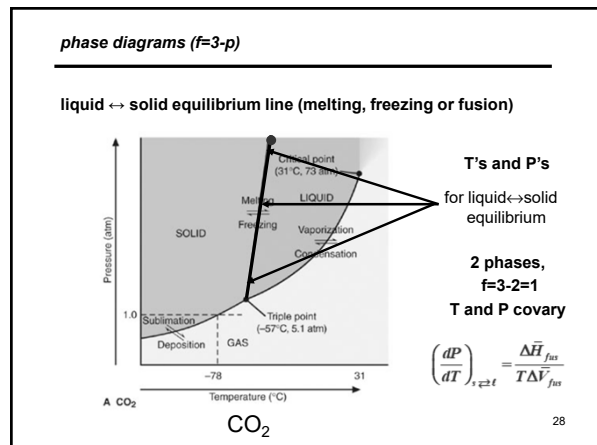
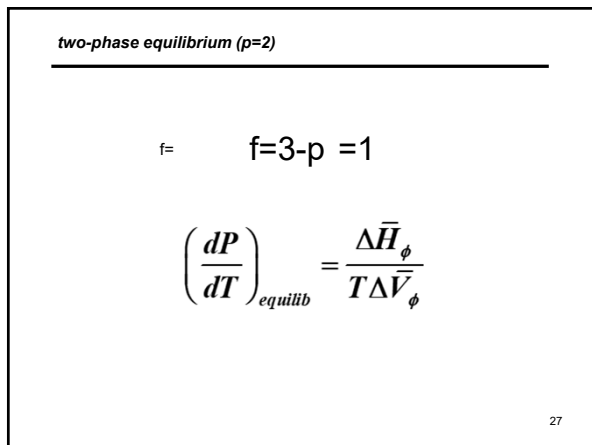
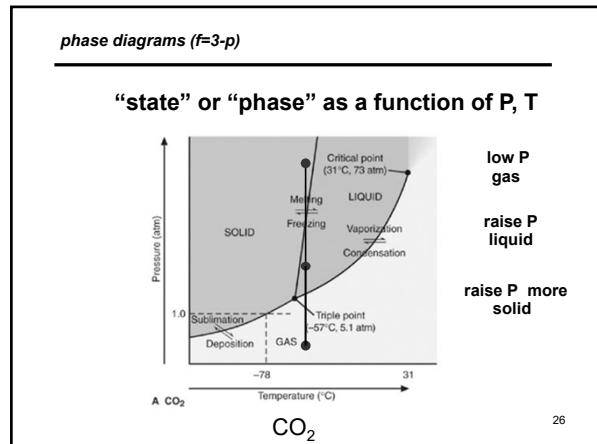
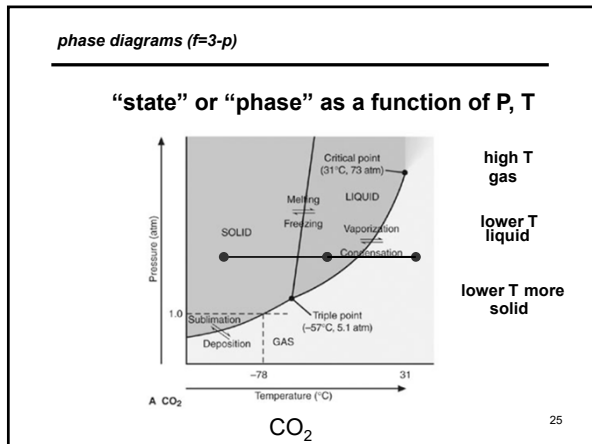
“state” or “phase” as a function of P, T



lower T more
 solid

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Lectures 18-19 Introduction to Phase Diagrams


critical point and triple point

- Triple point: for a pure substance, there is only one point (value of T and P) where all three phases (solid, liquid, and gas) can simultaneously exist in equilibrium
- Critical point: point (value of T and P) above which liquid and gas become one phase (fluid or supercritical fluid)



 movie: benzene critical point [A](#) [B](#)

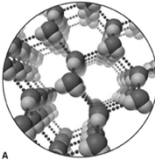

originally from: jchemed.chem.wisc.edu/jcesoft/ccca/samples/ccca2benzene.html

http://www.youtube.com/watch?v=79H2_OVBMGAA 

31

why does ice float ?

- H₂O is polar and can form hydrogen bonds (macho intermolecular forces)
- High surface tension and capillarity
- Hydrogen bonds form very open structure in solid H₂O (ice) giving ice a lower density than H₂O liquid. ICE FLOATS!!

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ice bomb !!!!





Originally at <http://www.jce.divched.org/JCESoft/CCA/irell/pages/cca2icebomb.html>

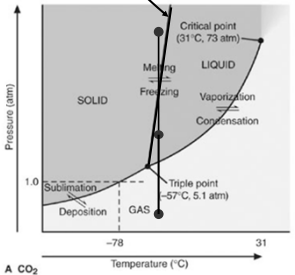
33

remember for CO₂ : P increases gas → liquid → solid

CO₂(s) ⇌ CO₂(l)

$$\left(\frac{dP}{dT}\right)_{j_{eq,t}} = \frac{\Delta\bar{H}_{melt}}{T\Delta\bar{V}_{melt}}$$

$\Delta\bar{H}_{melt} > 0;$
 $V_l > V_s; \Delta\bar{V} > 0$
 $\Rightarrow \left(\frac{dP}{dT}\right)_{j_{eq,t}} > 0$



low P
gas

raise P
liquid
(more dense)

raise P more
solid
(most dense)

A CO₂ CO₂

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phase diagram for water

remember for CO₂; P increases gas → liquid → solid

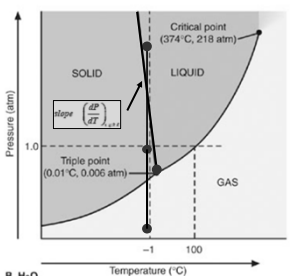
but for H₂O
as P increases:

gas
 ↓
 solid
 ↓
 liquid

WHY? $\left(\frac{dP}{dT}\right)_{j_{eq,t}} < 0$
for ice ⇌ water


$$\left(\frac{dP}{dT}\right)_{j_{eq,t}} = \frac{\Delta\bar{H}_{melt}}{T\Delta\bar{V}_{melt}}; \Delta\bar{H}_{melt} > 0$$

$V_l > V_s; \Delta\bar{V}_{melt} < 0 \Rightarrow \left(\frac{dP}{dT}\right)_{j_{eq,t}} < 0$



B H₂O



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ice skater myth  (ice skater real; Kristi Yamaguchi)

Does the weight of an ice skater create a pressure that melts ice to form a liquid groove for skate?

$$\ln \frac{T_{melt}}{T_{amb}} = \frac{(\bar{V}_l - \bar{V}_s)}{\Delta\bar{H}_{fusion}} (P - 1 \text{ atm})$$

$(\bar{V}_l - \bar{V}_s) < 0 \Rightarrow \text{pressure 'melts' ice}$

E&R^(4th) Problem P8.3
 508 bar for ΔT = -3.5°
 'thin blade'; 78kg = 172lb;
 P_{thin blade} = 243 bar; ΔT = -1.67°

NO, not even if they are quite 'weighty'!
 (not enough pressure and further details of water-ice phase diagram)

e.g. Rosenberg, Robert (December 2005). "Why is ice slippery?". *Physics Today*; 50-54.

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Lectures 18-19 Introduction to Phase Diagrams

Rosenberg, Robert (December 2005). "Why is ice slippery?". *Physics Today*: 50-54.

Figure 1. An ice skater exerts pressures on the order of a few hundred atmospheres on the ice surface, enough to reduce the melting temperature by only a few degrees. Premelting—the development of a liquid-like surface layer at temperatures below freezing—and frictional heating of the ice as skaters move around must account for ice's slipperiness at the wide variety of subzero temperatures found in nature. (Ice Skating, by Hy Sandham, 1885, courtesy of the Library of Congress)

The nature of the liquid-like layer is not clear from experimental measurements, so theorists have tried to clarify the situation.

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effects of inert gas (increased total pressure) on vapor pressure
E&R_{4th}, sec. 8.7

pure H₂O at 298K
 $P^*_{H_2O} = 0.0316 \text{ bar}$

1192751 ELIJI

'normal' H₂O in $P_{\text{total}} = 1 \text{ bar}$
[H₂O (g) + N₂ (g) + O₂ (g)] at 298K
 $P_{H_2O} = 0.031622 \text{ bar}$

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E&R_{4th} section 8.7 (effect of inert gas on vapor pressure)

Temperature = T

Pressure = $P_A + P_B$

Vapor phase

Condensed phase A

H₂O at 300 K
 $P^*_{H_2O} = 0.328 \text{ atm}$

add air (inert N₂ + O₂) to raise $P_{\text{total}} = 1 \text{ atm}$

new $P_{H_2O} = 0.32832 \text{ atm}$
 $(P/P^*)_{H_2O} = 1.00071$

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End of Lecture

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triple point

triple point: simultaneous equilibrium of gas, liquid solid

Pressure (atm)

Temperature (°C)

1.0

-78

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SOLID

LIQUID

GAS

Melting

Freezing

Vaporization

Condensation

Sublimation

Deposition

Critical point (31 °C, 73 atm)

Triple point (-57°C, 5.1 atm)

3 phases,
 $f = 3 - p$
 $f = 3 - 3 = 0$
T and P fixed

41

vary T and P through critical point

Pressure (atm)

Temperature (°C)

1.0

-78

31

SOLID

LIQUID

GAS

Melting

Freezing

Vaporization

Condensation

Sublimation

Deposition

Critical point (31 °C, 73 atm)

Triple point (-57°C, 5.1 atm)

fluid

42